

UNIVERSITY Faculty of M OF LJUBLJANA and Physics



## Generating Entangled Photon Pairs Using Nonlinear Crystal With Type-II SPDC

Seminar I - 1st year, 2nd cycle Ljubljana, Maj 2025

Author: Gaja Šalamun Advisor: dr. Peter Jeglič Co-advisor: dr. Izidor Benedičič



## Introduction

- $\ensuremath{\circ}$  reached field of application
- o promise:
  - more precise measurements
  - secure communication
  - implementation of otherwise overly complex simulations
- $_{\odot}$  scalling to larger systems
- o using photon polarization
- connect two stationary nodes





## **Quantum Network**

 $_{\odot}$  separate nodes with its own quantum system

- $\circ$  send light through optical fibers
  - state encoded in polarization, frequency or phase
  - fast speed of light
  - low losses
  - must be isolated
- secure data transmission
- o cannot amplify
- $_{\odot}$  limited to a few 100 km
- o entanglement swapping
- using satellites



## Entanglement

o unit is a qubit

- any general two-level quantum system
- $\circ$  single qubit state:  $|\psi\rangle = \sum_i c_i |\phi_i\rangle$
- $\circ$  two-qubit state:  $|\psi_{AB}\rangle = |\psi\rangle_A \otimes |\psi\rangle_B$
- $\circ$  photons  $\rightarrow$  polarization
  - $|\psi\rangle = \alpha |H\rangle + \beta e^{i\phi} |V\rangle$

o product state:

$$\begin{split} |\psi\rangle &= |\psi\rangle_1 \otimes |\psi\rangle_2 = \left(\alpha_1 |H\rangle_1 + \beta_1 e^{i\phi_1} |V\rangle_1\right) \otimes \left(\alpha_2 |H\rangle_2 + \beta_2 e^{i\phi_2} |V\rangle_2\right) = \\ &= c_1 |H\rangle_1 |H\rangle_2 + c_2 |H\rangle_1 |V\rangle_2 + c_3 |V\rangle_1 |H\rangle_2 + c_4 |V\rangle_1 |V\rangle_2 \end{split}$$

separable states

o statistical mixture of states

 $\Rightarrow$  density matrix formalism

#### Entanglement

 $\circ$  density operator:  $\hat{\rho} = \sum_{i} p_i |\psi_i\rangle \langle \psi_i |$ 

- pure states:  $\hat{\rho} = |\psi\rangle\langle\psi|$
- mixed states:  $\hat{\rho} = \sum_{i} p_i |\psi_i\rangle \langle \psi_i |$

 $\circ$  density matrix:

$$\rho = \frac{|HH\rangle}{|VV\rangle} \begin{pmatrix} A_1 & B_1 e^{i\phi_1} & B_2 e^{i\phi_2} & B_3 e^{i\phi_2} \\ B_1 e^{i\phi_1} & A_2 & B_4 e^{i\phi_4} & B_5 e^{i\phi_5} \\ B_2 e^{i\phi_2} & B_4 e^{i\phi_4} & A_3 & B_6 e^{i\phi_6} \\ B_3 e^{i\phi_3} & B_5 e^{i\phi_5} & B_6 e^{i\phi_6} & A_4 \end{pmatrix}$$

## Entanglement

A state is entangled if it is not separable, meaning we cannot write it as a product state.

- o measurements are correlated
- examples:
  - 1) four Bell states

$$\begin{split} \left|\psi^{\pm}\right\rangle &= \frac{1}{\sqrt{2}} \left(\left|H\right\rangle_{1}\left|V\right\rangle_{2} \pm \left|V\right\rangle_{1}\left|H\right\rangle_{2}\right) \\ \left|\phi^{\pm}\right\rangle &= \frac{1}{\sqrt{2}} \left(\left|H\right\rangle_{1}\left|H\right\rangle_{2} \pm \left|V\right\rangle_{1}\left|V\right\rangle_{2}\right) \end{split}$$

2) GHZ state

$$|GHZ\rangle = \frac{1}{\sqrt{2}} (|H\rangle^{\otimes n} \pm |V\rangle^{\otimes n})$$

## **Entanglement – Characterization**

quality of the entangled state

- $_{\odot}$  fidelity: how close the real state is to the targeted state
  - $F \sim \langle \rho_1 | \rho_2 \rangle$
  - $0 \leq F(\rho_1, \rho_2) \leq 1$

 $\circ$  purity: how much the state is mixed

- $\gamma = Tr\{\rho^2\}$
- $0 \le \gamma(\rho) \le 1$

o tangle: measure of entanglement

•  $0 \le T(\rho) \le 1$ 

## Spontaneous Parametric Downconversion – SPDC

- o second-order nonlinear optical effect
- strictly quantum effect
- splits into two photons
- o energy conservation
  - $\omega_p = \omega_s + \omega_i$
- momentum conservation
  - $\Delta \vec{k} = \vec{k}_p \vec{k}_s \vec{k}_i$
  - $\Rightarrow$  phase matching conditions
- inefficient process
  - ⇒ benefit









## SPDC – Nonlinear Crystals $k_p = k_s + k_i \xrightarrow{k_j = \frac{n_j \omega_j}{c}} \frac{n_p \omega_p}{c} = \frac{n_s \omega_s}{c} + \frac{n_i \omega_i}{c}$

 $\circ$  normal dispersion:  $n_i(\omega_i) < n_s(\omega_s) < n_p(\omega_p)$  for  $\omega_i < \omega_s < \omega_p$ 

 $\Rightarrow$  birefringent crystals

- $\circ$  refractive indices:  $n_e$  and  $n_o$
- o birefringent phase matching
  - $\Delta k = 0$
- o cone-like distribution



### SPDC – Walk-off

- $\circ$  input beam and generated beams do not travell in the same direction
- $\circ$  temporal walk-off:

$$\Delta t = \left| L_c \left( \frac{1}{n_e} - \frac{1}{n_o} \right) \right|$$

 $\circ$  spatial walk-off:

$$\delta = -\frac{1}{n_e} \frac{dn_e}{d\theta}$$

 $_{\odot}$  limits the length of the crystal



### SPDC – Quasi-phase Matching

 $\circ \Delta k = k_p - k_s - k_i \rightarrow 0$  ... hard to satisfy

 $\Rightarrow$  modify the material

o quasi-phase matching condition:

$$\vec{k}_p = \vec{k}_s + \vec{k}_i + \vec{K}_G$$
, where  $K_G = \frac{2\pi}{\Lambda}$ 

•  $\Lambda : 10 - 100 \, \mu m$ 

 $\circ$  no angle tuning

 $\circ$  collinear downconversion

○ travel along crystal axis

o can achieve type-0 SPDC



## SPDC – Periodic Poling

- $\circ$  2 methods:
  - 1) thin slices, where every other is rotated
  - 2) periodic poling
- o ferroelectric materials
  - LN: lithium niobate
  - LT: lithium tantalate
  - KTP: potassium titanyl phosphate linear susceptibility:  $\sim 10^{-11}$  m/V
- $\circ$  field: ~10 kV/mm
- need temperature tuning
  - n = n(T)



#### **SPDC** – Periodic Poling

 $\vec{k}_{p} = \vec{k}_{s} + \vec{k}_{i} + \vec{K}_{G} \xrightarrow{angle} k_{s} \cos \theta_{s} + k_{i} \cos \theta_{i} = k_{p} - K_{G}$  $k_{s} \sin \theta_{s} = k_{i} \sin \theta_{i}$ 

$$\circ$$
 can express:  $\vec{k}_j = \frac{2\pi \vec{n}_j}{\lambda_j}$ 

- o collinear type-II SPDC:
  - ellipsoid equation

$$n_s(\theta_s, T, \lambda_s) = \frac{n_x(T, \lambda_s) n_z(T, \lambda_s)}{\sqrt{n_x^2(T, \lambda_s) \cos \theta_s^2 + n_z^2(T, \lambda_s) \sin \theta_s^2}}$$

Sellmeier equations

$$n_p = n_y(T, \lambda_p)$$
$$n_i = n_y(T, \lambda_i)$$



# Generation of Polarized Entangled Photon Pairs

1) single emitter for type-II SPDC



2) sandwich scheme for type-I SPDC



- 3) interferometric methods
  - Mach-Zehnder interferometer
  - Sagnac interferometer

## Sagnac Interferometer

- o phase-stable
- $\circ$  QPM:  $|H_p\rangle \rightarrow |H_s\rangle, |V_i\rangle$ 
  - $\Rightarrow$  helps with imperfections of PBS
- o pump polarization
- $\circ$  finer tuning
- plano-convex lens
- o dichroic mirror
- band pass filter and multiplexer



#### Sagnac Interferometer – loop $|H\rangle$



#### Sagnac Interferometer – loop $|V\rangle$





# Entanglement Measurement and Its Interpretion

- quantum state tomography
  - maximum likelihood method
- $\circ$  reconstruct the density matrix
- o external setup
- compensating wave plates
- choose measurement basis
- project in different bases
  - linear
  - diagonal
  - circular



#### **Tomography Results**

1) density matrix

$$\rho = \frac{|HH\rangle}{|VH\rangle} \begin{pmatrix} A_1 & B_1 e^{i\phi_1} & B_2 e^{i\phi_2} & B_3 e^{i\phi_2} \\ B_1 e^{i\phi_1} & A_2 & B_4 e^{i\phi_4} & B_5 e^{i\phi_5} \\ B_2 e^{i\phi_2} & B_4 e^{i\phi_4} & A_3 & B_6 e^{i\phi_6} \\ B_3 e^{i\phi_3} & B_5 e^{i\phi_5} & B_6 e^{i\phi_6} & A_4 \end{pmatrix}$$

2) histogram for real and imaginary amplitudes



#### Implementation of Experiment – IJS







## Implementation of Experiment – IJS

 $\circ$  used:

- 780 nm pump laser
- PPLN crystal ( $50 \times 2 \times 1 \text{ mm}$ )
- 1550 nm downconverted photons
- 250 mm plano-convex lens

 $\circ$  results:

- fidelity: (97.1  $\pm$  0.08) %
- purity: (97.9 ± 0.09) %
- tangle: (92.5 ± 0.07) %







#### Implementation of Experiment Over Record Long Distance

- Austrian Academy of Science
- $\circ$  over 248 km in optical fibers
- $\circ$  used:
  - 775.06 nm pump laser
  - PPLN crystal
  - 1550.12 nm downconverted photons
- $\circ$  fidelity: > 99%
- independent of weather



## Conclusion

○ basis for:

- realization of quantum communication
- quantum teleportation
- o important for quantum key distribution (QKD)
- $_{\odot}$  transfer of information between particles that are far apart
- $_{\odot}$  robust and easy to use method

